

Chapter "4"

Alpha Decay

Until the discovery of spontaneous fission, α - **decay** was the only known type of radioactive disintegration, in which nuclei emits heavy charged particles (${}^4_2\text{He}_2$), i.e. α - **particle** is the nucleus of the Helium atom. Alpha particle is a very **stable** and tightly bound structure. Alpha emission is a coulomb repulsion effect, it becomes increasingly important for heavy nuclei because the disruptive coulomb force increases as (Z^2) faster than the nuclear binding force (**which increases approximately as A**). Most nuclei with $A > 190$ and (many with $150 < A < 190$) are energetically can decay by α - **emission**. In α - **decay** the nuclear reaction can be written in this form :



All the α - **ray** emitted by heavy nuclei ($A > 190$) can be classified into one of the **four** independent decay chains with mass numbers $4n$, $4n + 1$, $4n + 2$, where (n) is an integer. The decay process will tend to concentrate the nuclei in the longest - lived member of the chain.

① The Thorium ($4n$) chain :

Every member of this chain has a mass number which is a multiple of **four**, i.e. $A = 4n$. The longest - lived member ${}^{232}_{90}\text{Th}$ has a half - life, $T_{1/2} = 1.41 \times 10^{10} \text{ year} = 14.1 \times 10^9 \text{ year}$, which is about **Three times** the age of the earth ($T_{\text{earth}} = 4.5 \times 10^9 \text{ year}$).

② The Neptunium ($4n + 1$) chain :

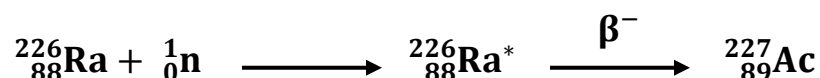
In nature there is no $A = 4n + 1$ series, This is because its longest lived member ${}^{237}_{93}\text{Np}$ has a half - life of only $2.14 \times 10^6 \text{ year}$. This series can be produced artificially. The ($4n + 1$) series terminates at an stable ${}^{209}_{83}\text{Bi}$, this is the only series which do not terminate at an isotope of **Pb**.

③ The Uranium ($4n + 2$) chain :

This is the longest known series. In nature it begins with the heaviest naturally occurring nucleus ${}^{238}_{92}\text{U}$ ($T_{1/2} = 4.47 \times 10^9 \text{ year}$), which is the age of the earth.

④ The Actinium ($4n + 3$) chain :

This is the only naturally occurring nuclei which undergoes fission by slow neutrons . **Actinium** $^{227}_{89}\text{Ac}$ can be separated from **Uranium** minerals . The β – ray spectrum of **actinium** has a maximum energy of only **0.04 MeV** . **Actinium** can be produced artificially in reactors by irradiation of $^{226}_{88}\text{Ra}$, the nuclear reaction is :



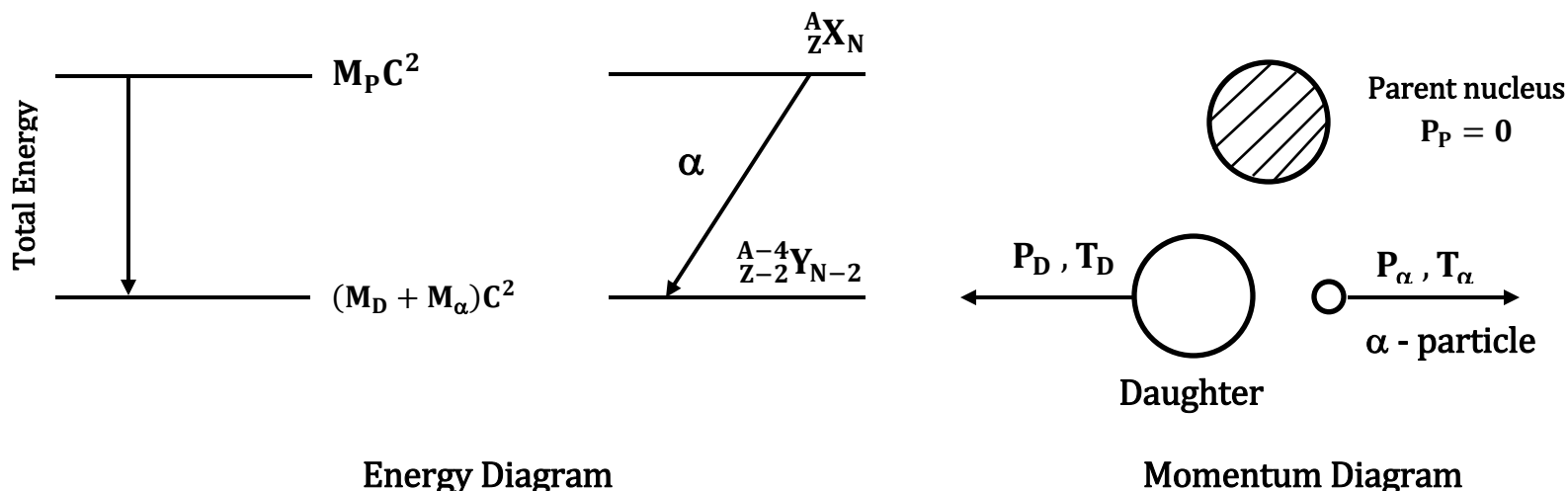
Fissionable materials ^{235}U and ^{239}Pu are belong to this series.

Table : Some characteristics of the Disintegration series of the Heavy Elements

Name of series	Type	Final Nucleus (Stable)	<u>Longest Lives Member</u>	
			Nucleus	Half – life (years)
Thorium	$4n$	^{208}Pb	^{232}Th	14.1×10^9
Neptunium	$4n + 1$	^{209}Bi	^{237}Np	2.14×10^6
Uranium	$4n + 2$	^{206}Pb	^{238}U	4.47×10^9
Actinium	$4n + 3$	^{207}Pb	^{235}U	7.04×10^8

Energetics of α - decay :

The spontaneous emission of α - **particle** can be represented as follows :



Energy Diagram

Momentum Diagram

In α - **decay** we assume the initial decaying nucleus (**parent nucleus**) to be at rest ($T_P = 0, P_P = 0$), then the energy of the initial system is just the rest energy of the parent nucleus $M_P C^2$. The final state (**after decay**) consist of daughter and α - **particle**, from the conservation of **energy** and **momentum** we have :

$$M_P C^2 = M_D C^2 + T_D + M_\alpha C^2 + T_\alpha \quad \text{-----} \quad (1)$$

$$0 = \vec{P}_D + \vec{P}_\alpha \quad \text{-----} \quad (2)$$

Where T_α and \vec{P}_α are the kinetic energy and momentum of α - **particle**

T_D and \vec{P}_D are the kinetic energy and momentum of **daughter nucleus**

$$\therefore M_P C^2 = (M_D + M_\alpha) C^2 + (T_D + T_\alpha) \quad \text{-----} \quad (3)$$

The decay energy is the net energy released in the decay called the **Q - value** is the sum of the resultant kinetic energies, i.e.

$$Q = T_D + T_\alpha \quad \text{-----} \quad (4)$$

$$\text{Therefore} \quad M_P C^2 = (M_D + M_\alpha) C^2 + Q$$

$$\therefore Q = [M_P - (M_D + M_\alpha)] C^2 \quad \text{-----} \quad (5)$$

$$Q_\alpha = (M_P - M_D + M_\alpha) C^2 \quad \text{-----} \quad (6)$$

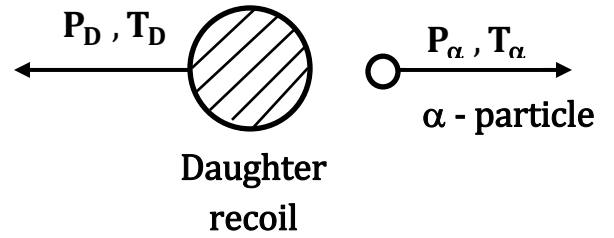
And the decay occurs spontaneously if $Q > 0$, the **Q - value** can be calculated from atomic masses given in tables (**in amu**) then multiplying by 931.5 (1 amu = 931.5 MeV) to have the **Q - value** in MeV.

Since the **Q - value** is given by eq. (4) , $Q = T_D + T_\alpha$ and since the kinetic energy $T \ll mc^2$, we can safely use nonrelativistic expressions to evaluate T_D and T_α .

From eq. (2) we have $P_D = -P_\alpha$

$$T_D = \frac{P_D^2}{2M_D} = \frac{P_\alpha^2}{2M_D} = \frac{Z M_\alpha T_\alpha}{Z M_D} = \frac{M_\alpha}{M_D} T_\alpha$$

$$\therefore T_D = \frac{M_\alpha}{M_D} T_\alpha \quad \text{-----} \quad (7)$$



The recoil kinetic energy of daughter T_D is not negligible as it is in the case of γ - decay because α - particle is a large mass.

$$\text{Since } Q_\alpha = T_D + T_\alpha = \frac{M_\alpha}{M_D} T_\alpha + T_\alpha = \frac{M_D + M_\alpha}{M_D} T_\alpha$$

Q_α طاقة انحلال الفا

T_α طاقة حركية لألفا

$$Q_\alpha = \frac{M_P}{M_D} T_\alpha = \left(\frac{M_P}{M_P - 4} \right) T_\alpha$$

and since $M_P \approx A$; (A is the mass number of parent nucleus)

$$\therefore Q_\alpha \cong \left(\frac{A}{A - 4} \right) T_\alpha \quad \text{-----} \quad (8)$$

$$\text{or } T_\alpha = Q_\alpha \left(\frac{A}{A - 4} \right) = Q_\alpha \left(1 - \frac{4}{A} \right)$$

$$T_\alpha = Q_\alpha \left(1 - \frac{4}{A} \right) \quad \text{-----} \quad (9)$$

Alpha decay cannot take place unless $Q_\alpha > 0$, i.e. is positive and the separation energy of α - particle S_α is equal to :

$$Q_\alpha = -S_\alpha$$

$$Q_\alpha = B_{\text{tot.}} \left({}_{Z-2}^{A-4}\text{X}_{N-2} \right) + B_{\text{tot.}}(\alpha) - B_{\text{tot.}} \left({}_Z^A\text{X}_N \right)$$