

College of Science
Department of Physics
Fourth Class
Lecture 13

Quantum Mechanics

2023-2024

Lecture 13: Free particle in a box

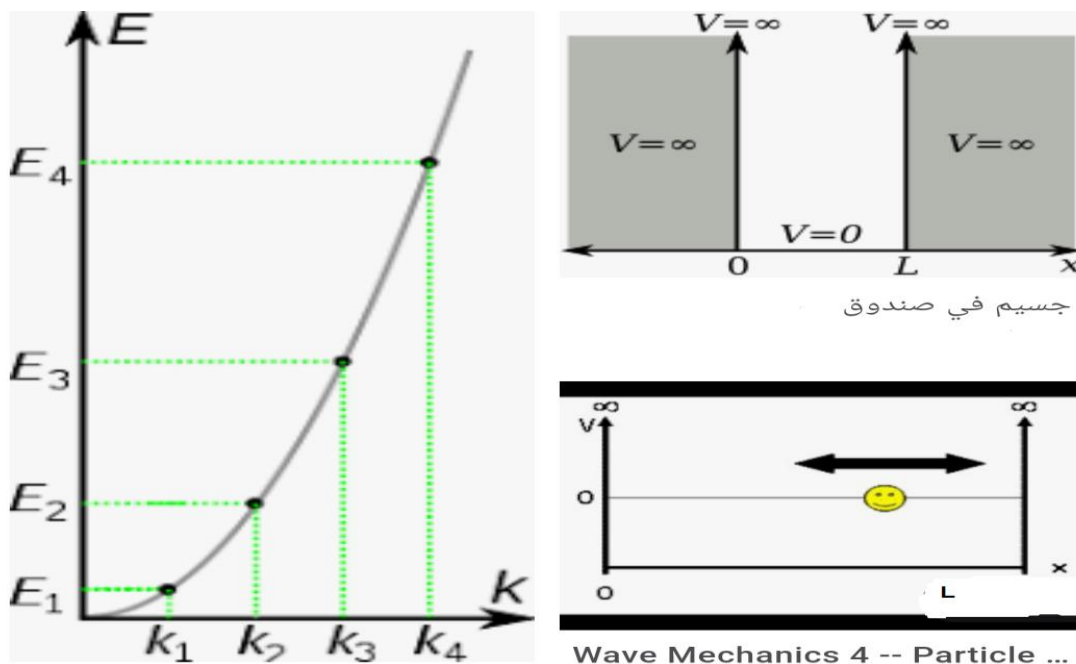
Preparation

Assist. prof. Alaa Abdul Hakeim

6.2 Free particle in a box

Let a single particle of mass "m" be confined within a rectangular potential box with edges parallel to X,Y & Z – axes of length l_x, l_y, l_z respectively. The particle can move freely within the region

$$0 < x < l_x \quad , \quad 0 < y < l_y \quad , \quad 0 < z < l_z$$



i.e. inside the box, the potential function $V(x,y,z)$ is equal zero, but it is ∞ ever where outside the box. Therefore, the particle will rebound when it will strike either of the boundaries.

The schrödinger eq. for the particle is

$$-\frac{\hbar^2}{2m} \nabla^2 \Psi = E \Psi \quad \rightarrow V = 0$$

$$\text{Or } \nabla^2 \Psi + \frac{2m}{\hbar^2} E \Psi = 0 \quad \rightarrow V = 0$$

E is total energy of the particle

m is mass of the particle

$$\frac{\partial^2 \Psi}{\partial x^2} + \frac{\partial^2 \Psi}{\partial y^2} + \frac{\partial^2 \Psi}{\partial z^2} + \frac{2m}{\hbar^2} E \Psi = 0 \text{ ---8}$$

This eq. solved by separation of variables.

$$\Psi(x,y,z)= X(x) Y(y) Z(z)$$

Substitute in eq.8 and divide by Ψ we get

$$\frac{1}{X} \frac{\partial^2 X(x)}{\partial x^2} + \frac{1}{Y} \frac{\partial^2 Y(y)}{\partial y^2} + \frac{1}{Z} \frac{\partial^2 Z(z)}{\partial z^2} + \frac{2m}{\hbar^2} E = 0$$

$$\frac{1}{X} \frac{\partial^2 X(x)}{\partial x^2} = -\frac{1}{Y} \frac{\partial^2 Y(y)}{\partial y^2} - \frac{1}{Z} \frac{\partial^2 Z(z)}{\partial z^2} - \frac{2m}{\hbar^2} E$$

$$\text{Let } \frac{1}{X} \frac{\partial^2 X(x)}{\partial x^2} = k_x^2 \quad k_x - \text{const.} \quad - - - - - 9$$

$$k_x^2 = -\frac{1}{Y} \frac{\partial^2 Y(y)}{\partial y^2} - \frac{1}{Z} \frac{\partial^2 Z(z)}{\partial z^2} - \frac{2m}{\hbar^2} E$$

$$\frac{1}{Y} \frac{\partial^2 Y(y)}{\partial y^2} = -k_x^2 - \frac{1}{Z} \frac{\partial^2 Z(z)}{\partial z^2} - \frac{2m}{\hbar^2} E$$

$$\text{Let } \frac{1}{Y} \frac{\partial^2 Y(y)}{\partial y^2} = k_y^2 \quad \text{-----} 10$$

$$-k_x^2 - \frac{1}{Z} \frac{\partial^2 Z(z)}{\partial z^2} - \frac{2m}{\hbar^2} E = k_y^2$$

This eq. may be written as

$$\frac{1}{Z} \frac{\partial^2 Z(z)}{\partial z^2} = k_z^2 \quad - - - - - 11$$

$$= -k_x^2 - k_y^2 - \frac{2m}{\hbar^2} E$$

$$k_x^2 + k_y^2 + k_z^2 = -\frac{2m}{\hbar^2} E$$

$$k_x^2 = -\frac{2m}{\hbar^2} E_x \quad , \quad k_y^2 = -\frac{2m}{\hbar^2} E_y \quad , \quad k_z^2 = -\frac{2m}{\hbar^2} E_z$$

The eqs.(9,10,11) may be written as

$$\frac{\partial^2 X(x)}{\partial x^2} + \frac{2m}{\hbar^2} E_x X = 0 \quad - - - - a$$

$$\frac{\partial^2 Y(y)}{\partial y^2} + \frac{2m}{\hbar^2} E_y Y = 0 \quad - - - - b$$

&

$$\frac{\partial^2 Z(z)}{\partial z^2} + \frac{2m}{\hbar^2} E_z Z = 0 \quad \text{--- --- ---} \quad 12$$

The general solution of eq. 12a will be sine function of arbitrary amplitude, frequency, and phase, i.e.

$$X(x) = A \sin(Bx + C) \quad \text{-----} \quad 13$$

Where A, B & C are constants to be determined from consideration of the boundary conditions as

$|\Psi|^2$ represent the probability of finding the particle at any point within the box.

$|X(x)|^2$ is the probability of finding the particle at any point along x-axis

As the pot. $\rightarrow \infty$ "Very high" at the walls of the box, the probability of finding the particle at walls will be zero, i.e.

$$|X(x)|^2 = 0 \quad \text{when } x=0 \quad \& \quad x=l_x$$

$$\text{Or } X(x)=0 \quad \text{when } x=0 \quad \& \quad x=l_x$$

Using these boundary condition, eq.13 gives

$$X(x) = A \sin(Bx + C) = 0$$

$$\sin C=0 \quad \text{when } x=0 \quad \text{and } \sin(Bl_x + C)=0 \quad \text{when } x=l_x$$

$$\text{which yield } C=0 \quad \& \quad \sin Bl_x = 0$$

$$\text{or } Bl_x = n_x \pi \quad n_x = 1, 2, 3, \text{-----+ve integer}$$

$$B = \frac{n_x \pi}{l_x}$$

Substituting values of B & C in eq. 13

$$X(x) = A \sin \frac{n_x \pi x}{l_x} \quad \text{-----} \quad 14$$

For simplicity we assume A to be a real number, then using the condition of normalize

$$\int_0^{l_x} |X(x)|^2 dx = 1$$

We have $\int_0^{l_x} \left| A \sin \frac{n_x \pi x}{l_x} \right|^2 dx = 1$

Or $A^2 \int_0^{l_x} \sin^2 \frac{n_x \pi x}{l_x} dx = 1$

$$A^2 \frac{l_x}{2} = 1 \quad \rightarrow \quad A = \sqrt{\frac{2}{l_x}}$$

Substituting the value of A in eq.14, the normalized function X(x) is given by

$$X(x) = \sqrt{\frac{2}{l_x}} \sin \frac{n_x \pi x}{l_x} \text{----- 15}$$

$$\frac{\partial^2 X}{\partial x^2} = - \left(\frac{n_x \pi}{l_x} \right)^2 X(x) \text{----- 16}$$

Sub.15 & 16 in to eq.12a , we get

$$\begin{aligned} \frac{\partial^2 X(x)}{\partial x^2} + \frac{2m}{\hbar^2} E_x X &= 0 \\ - \left(\frac{n_x \pi}{l_x} \right)^2 X(x) + \frac{2m}{\hbar^2} E_x X(x) &= 0 \end{aligned}$$

Or $\boxed{E_x = \frac{1}{2m} \left(\frac{n_x \pi \hbar}{l_x} \right)^2 = \frac{n_x^2 \hbar^2}{8ml_x^2}} \text{---17} \quad \text{when } \hbar = \frac{h}{2\pi}, n=1,2,3 \text{ ---}$

Note :1- energy is quantized E1, E2,E3,-----

2- $E_{min} = E_1 \neq 0$

Similarly, we may solve eqs.12b &12c and obtain

$$Y(y) = \sqrt{\frac{2}{l_y}} \sin \frac{n_y \pi y}{l_y} \quad , n_y = 1,2,3, \text{-----}$$

$$Z(z) = \sqrt{\frac{2}{l_z}} \sin \frac{n_z \pi z}{l_z} \quad , \quad n_z = 1, 2, 3, \dots$$

$$\boxed{E_y = \frac{n_y^2 h^2}{8m l_y^2}} \dots\dots\dots 18 \quad \boxed{E_z = \frac{n_z^2 h^2}{8m l_z^2}} \dots\dots\dots 19$$

Having

$$k_x^2 + k_y^2 + k_z^2 = -\frac{2m}{\hbar^2} E$$

And

$$k_x^2 = -\frac{2m}{\hbar^2} E_x \quad , \quad k_y^2 = -\frac{2m}{\hbar^2} E_y \quad , \quad k_z^2 = -\frac{2m}{\hbar^2} E_z$$

$$E = E_x + E_y + E_z$$

Eqs. 17,18,19 give the allowed values of energy levels along X, Y & Z axes respectively.

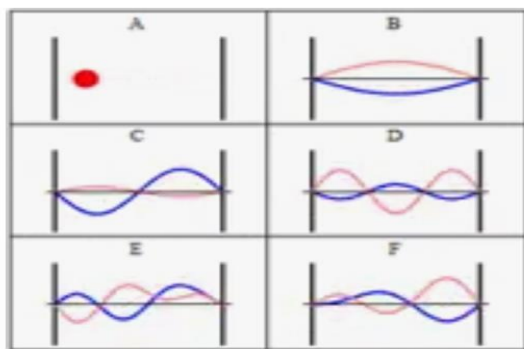
The allowed values of total energy are given by

$$E = E_x + E_y + E_z = \frac{h^2}{8m} \left[\frac{n_x^2}{l_x^2} + \frac{n_y^2}{l_y^2} + \frac{n_z^2}{l_z^2} \right]$$

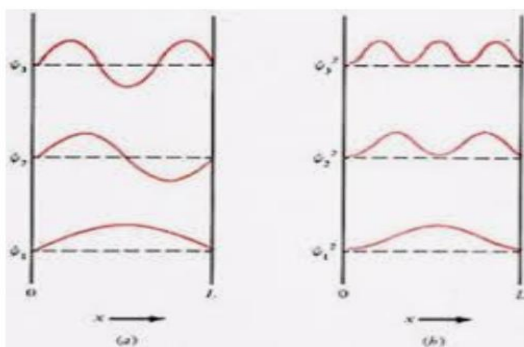
Where n_x, n_y, n_z denote any set of three positive numbers.

The complete wave – function Ψ_{n_x, n_y, n_z} has the form, for various values of the values of the quantum numbers n_x, n_y & n_z .

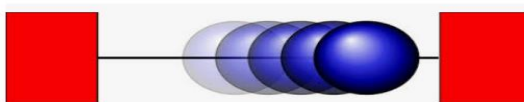
$$\begin{aligned} \Psi_{n_x, n_y, n_z}(x, y, z) &= X(x)Y(y)Z(z) \\ &= 2 \sqrt{\frac{2}{l_x l_y l_z}} \sin \frac{n_x \pi x}{l_x} \sin \frac{n_y \pi y}{l_y} \sin \frac{n_z \pi z}{l_z} \end{aligned}$$



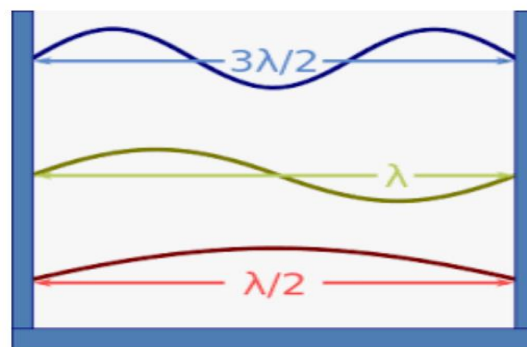
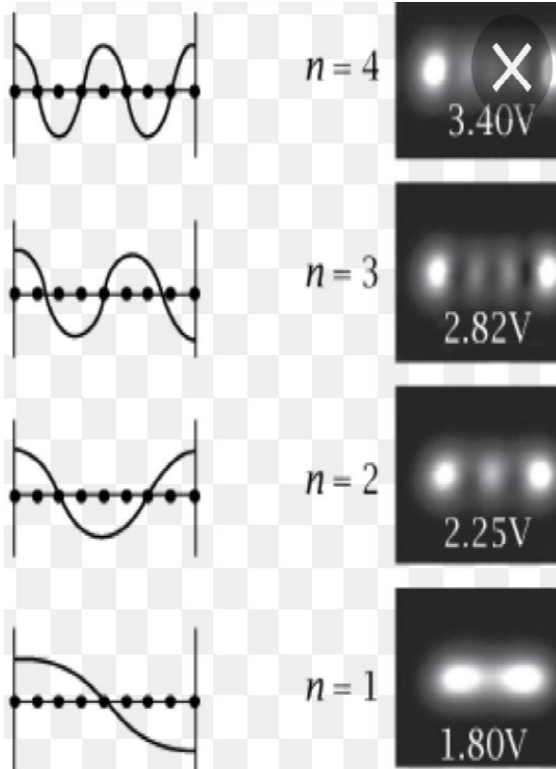
جسيم في صندوق - المعرفة
marefa.org



محاضرة 12 فيزياء حديثة - شبكة الف...
hazemsakeek.net



جسيم في صندوق - Wikiwand
wikiwand.com



Problem 1:

Prove that the wave functions of a particle in one dimensional box ($l_x = a$) are orthogonal.

Solution:

$$\Psi(x) = \sqrt{\frac{2}{l_x}} \sin \frac{n_x \pi x}{l_x}$$

$$\Psi_n(x) = \sqrt{\frac{2}{a}} \sin \frac{n_x \pi x}{a} \quad \text{when } l_x = a$$

The orthogonal condition is

$$\int_{-\infty}^{\infty} \Psi_n^* \Psi_l dx = 0 \quad n \neq l$$

In free particle wave function in a box

$$\begin{aligned} \frac{2}{a} \int_0^a \sin \frac{n\pi x}{a} \sin \frac{l\pi x}{a} dx &= \frac{1}{a} \int_0^a 2 \sin \frac{n\pi x}{a} \sin \frac{l\pi x}{a} dx \\ &= \frac{1}{a} \int_0^a \left(\cos \frac{(n-l)\pi x}{a} - \cos \frac{(n+l)\pi x}{a} \right) dx \\ &= \frac{1}{a} \left[\frac{a}{(n-l)\pi} \sin \frac{(n-l)\pi x}{a} - \frac{a}{(n+l)\pi} \sin \frac{(n+l)\pi x}{a} \right]_0^a = 0 \end{aligned}$$

Because $n-l \rightarrow \text{integer}$ & $n+l \rightarrow \text{integer}$

So $\sin 0 = 0$

$$\sin(n-l)\pi = 0 \quad \& \quad \sin(n+l)\pi = 0$$

$$\frac{2}{a} \int_0^a \sin \frac{n\pi x}{a} \sin \frac{l\pi x}{a} dx = \text{zero}$$

The wave functions are orthogonal

Problem 2:

a- Verify that a travelling wave solution of the schrödinger eq. for a free particle

$\Psi(x,t) = A \exp i(k_x x - \omega t)$ is also an eigen function of linear momentum

b- Evaluate the probability density for the same travelling wave solution and show that both results are consistent with Heisenberg uncertainty principle.

Solution:

$$\begin{aligned} \text{a- } \Psi(x, t) &= A \exp i(k_x x - \omega t) \\ &= A \exp(i k_x x) \exp\left(-i \frac{Et}{\hbar}\right) \end{aligned}$$

Where E is the energy eigen value and

$$k_x = \frac{\sqrt{2mE_x}}{\hbar}$$

$$\begin{aligned} \hat{P}\Psi(x,t) &= \frac{\hbar}{i} \frac{\partial}{\partial x} \Psi(x, t) = \frac{\hbar}{i} \frac{\partial}{\partial x} A \exp(i k_x x) \exp\left(\frac{-iEt}{\hbar}\right) \\ &= \hbar k_x \Psi(x, t) \end{aligned}$$

$\hbar k_x$ is eigen value of linear momentum $P = \hbar k_x = \sqrt{2mE_x}$

thus, for a fixed energy the linear momentum is uniquely defined.

b- The prob.density = $\Psi^*(x, t)\Psi(x, t) = A^* A$

i.e. This is independent of X and thus the X-coordinate is completely unknown

i.e. As the momentum is uniquely defined, the position cannot be known ($\Delta P = 0$, $\Delta x = \infty$).

Problem: H.W

- 1- A crystalline chip in a cell phone has a side length of (0.2nm). Calculate the minimum speed of an electron move inside it.
- 2- Calculate the maximum speed of a particle has mass (1gm) inside the box a side length (1 cm). calculate the A- particle's energy if the velocity is (3 cm/s). B- n to that energy level.